# NATIONAL ADVISORY COMMITTEE FOR AERONAUTICS

TECHNICAL NOTE

No. 1187

FAIRCHILD AND THEFT FILDPEKLY
ENGINEERING DEPARTMENT

FORMULAS FOR ADDITIONAL MASS CORRECTIONS

TO THE MOMENTS OF INERTIA OF AIRPLANES

By Frank S. Malvestuto, Jr. and Lawrence J. Gale

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TO THE MOMENTS OF INERTIA OF AIRPIANES

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#### SUMMARY

Formulas are presented for the calculation of the additional-mass corrections to the moments of inertia of airplanes. These formulas are of particular value in converting the virtual moments of inertia of airplanes or models experimentally determined in air to the true moments of inertia. A correlation of additional moments of inertia calculated by these formulas with experimental additional moments of inertia obtained from vacuum—chamber tests of 40 spin—tunnel models indicates that the formulas give satisfactory estimations of the additional moments of inertia.

#### INTRODUCTION

In stability investigations involving free-flight tests of dynamically scaled models or of full-scale airplanes it is necessary to know accurately the true moments of inertia of the model or airplane. In order to obtain the true moments of inertia of the model or airplane the experimental moments of inertia, determined by the pendulum method, must be corrected for the effect of the surrounding air. The effect of this ambient air on the apparent moments of inertia is usually small but may be as large as 25 percent of the true moments of inertia for airplanes with low wing loading.

The fundamental basis of the effect of the ambient-air mass on bodies undergoing acceleration has been developed in references 1 and 2. References 3, 4, and 5 present experimental data on the effect of the ambient-air mass on the moments of inertia of flat plates obtained by swinging the plates as an integral part of a pendulum. Reference 3 presents a method of correcting the moments of inertia of a full-scale airplane for the ambient-air mass effect by considering the projected areas of various components of the airplane in planes normal and parallel to the plane of symmetry and

applying to these projected areas the experimentally determined additional moment-of-inertia corrections for flat plates of finite aspect ratio.

In the present paper, formulas are presented for the rapid evaluation of the additional-mass and moment-of-inertia corrections for airplanes that are swing as an integral part of a pendulum. The method proposed in the preceding reports has generally been conformed to with the exception that the ambient-dir mass effect for parts of the airplane such as the fuselage was determined theoretically from the motion of an ellipsoid in a three-dimensional potential flow. The method presented has been applied to 40 complete dynamic airplane models previously tested and the results have been compared with the experimentally determined values.

#### SYMBOIS

For convenience in defining certain symbols, sketches for identifying dimensional parts of the airplane are given in figures 1 and 2. The numerical values given in these figures are for use in an illustrative example that is subsequently presented.

```
span of surface (includes span of fusclage between surface
ъ
              panels; for vertical tail see fig. (b))
          area of surface (for wing and horizontal tail includes
S
              area of fuselage between surface panels)
          aspect ratio of surface (b2/S)
Α
          root chord of surface
c_r
          tip chord of surface \left(\frac{2S}{b} - c_r\right)
C+
c
          mean chord of surface (S/t)
          plan-form taper ratio of surface (c_r/c_+)
λ
           dihedral of wing or horizontal tail in degrees
           length of fuselage (see fig. 1)
L_{\Gamma}
           geometric average width of fuselage
           geometric average depth of fuselage
đ
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component in plane of surface of perpendicular distance

between axis of rotation and centroid of area of surface

- distance from centroid of side area of fuselage to axis of rotation parallel to and in the plane of the X-axis (conveniently referred to herein as axis of X swinging)
- component of distance in the X-Y principal plane of fuselage of the perpendicular distance between the centroid of plan area of fuselage and the axis of rotation parallel to and in the plane of the Y-axis
- distance from centroid of side area of fuselage to axis of rotation parallel to end in the plane of the Z-axis
- distance from centroid of vertical-tail area to axis of rotation parallel to and in the plane of the X-axis
- component of distance in the X-Y plane of fuselage of the perpondicular distance between the centroid of herizontal-tail area and the axis of rotation parallel to and in the plane of the Y-axis
- distance from centroid of vertical-tail area to axis of rotation for Z swinging
- lwy component of distance in the X-Y plane from the centroid of area of wing to axis of rotation parallel to and in the plane of the Y-axis
- k coefficient of additional mass of a flat rectangular plate
- k\* coefficient of additional moment of inertia of a flat rectangular plate
- D taper-ratio correction factor
- D<sub>P</sub> dihedral correction factor
- k<sub>fX</sub>, k<sub>fY</sub>, k<sub>fZ</sub> coefficients of additional mass of equivalent ellipsoids for motion along the X-, Y-, and Z-axes, respectively
- k: fX, k'fY, k'fZ coefficients of additional moments of inertia of equivalent ellipsoids about the X-, Y-, and Z-axes, respectively
- m<sub>a</sub> additional mass of a body

I<sub>a</sub> additional moment of inertia of a body

 $I_X$ ,  $I_Y$ ,  $I_Z$  moments of inertia about X-, Y-, and Z-body axes, respectively

 $I_{X_a}$ ,  $I_{Y_a}$ ,  $I_{Z_a}$  total additional moments of inertia about X-, Y-, and Z-body axes, respectively

IXa, IYa, IZa, total additional moments of inertia about X, Y, and Z swinging axes, respectively

density of air, slug per cubic foot

Subscripts:

W wing

fus fuselage

ht horizontal tail

vt vertical tail

#### DEVELOPMENT OF EQUATIONS

Wings and Tail Surfaces

In order to evaluate the additional-mass corrections for the wings and tail surfaces of an airplane it was assumed that these surfaces were flat plates and that the additional-mass corrections previously obtained for flat plates (references 5 and 6) were approximately correct for the wings and tail surfaces. The additional mass of a wing or tail surface in translatory motion may be determined from the following equation:

$$\mathbf{m}_{\mathrm{R}} = \frac{\pi \rho}{h} \, \mathrm{kc}^2 \mathrm{b} \tag{1}$$

where k, the coefficient of additional mass plotted in figure 3 as a function of the aspect ratio, has been obtained by averaging the experimental results of NACA tests presented in reference 5. The values presented in references 3 and 5 for a rectangular wing and those of the present analysis are assumed to be accurate for a tapered wing when considered as an equivalent rectangular wing of chord  $\bar{c}$ .

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The additional moment of inertia of a wing or tail surface rotating about its chord at the midspan is determined from the equation

$$I_{a} = \frac{\pi \rho}{48} k^{\epsilon} \bar{c}^{2} b^{3} \tag{2}$$

where k\*, the coefficient of additional moment of inertia plotted in figure 4 as a function of the aspect ratio of the plate, has been obtained by averaging the experimental results of NACA tests presented in reference 5. Application of correction factors for the effect of dihedral angle and taper ratio gives the following equation:

$$I_{a} = \frac{\pi \rho}{48} D_{\lambda} D_{\Gamma} k^{\dagger} \bar{c}^{2} b^{3}$$
 (3)

The correction factors for the effect of dihedral angle and taper ratio on the values of the additional mements of inertia for flat plates as given in reference 5 are presented in figures 5 and 6.

For rotation about a spanwise axis through the centreid of area of the surface the additional moment of inertia is given by

$$I_{e} = \frac{\pi \rho}{48} k^{\dagger} \bar{c}^{3} b^{2} \tag{4}$$

where the aspect ratio used in determining  $k^2$  from figure 4 is now 1/A.

For rotation about an axis displaced from the centroid of area of the surface, the additional moment of inertia about the axis of rotation may be found from the expression

$$I_a^* = I_a + m_a l^2$$

where 1 is the component in the plane of the surface of the perpendicular distance between the axis of rotation and centroid of area of the surface. For example, the additional moment of inertia about an axis parallel to the chord of the surface is given by

$$I_{a}^{\dagger} = \frac{\pi \rho}{48} k^{2} \bar{c}^{2} b^{3} + \frac{\pi \rho}{4} k \bar{c}^{2} b l^{2}$$
 (5)

#### Fuselage

In order to evaluate the additional-mass and moment-of-inertia corrections for the fuselage it may be assumed that the fuselage of an airplane can be approximated in shape by an ellipsoid and that the values of additional mass and additional moments of inertia calculated for the ellipsoid are approximately correct for a fuselage which has the same length and volume as the ellipsoid. That is, the maximum depth and the maximum width of the equivalent

ellipsoid are equal, respectively, to  $\sqrt{\frac{6}{\pi}} \, d$  and  $\sqrt{\frac{6}{\pi}} \, w$  of the fuselage where d and w are values of the average depth and width of the fuselage.

For a fuselage moving in a direction parallel to one of its principal axes in an inviscid fluid, the additional mass is determined from the equations of linear momentum for the equivalent ellipsoid presented in reference 2. For motion along the Y and Z principal axes, the values of the additional mass in terms of the average depth and width of the fuselage are given, respectively, by

$$m_{a} = \rho k_{fY} L_{f} wd$$
 (6)

and

$$m_{a} = \rho k_{17} L_{f} wd \tag{7}$$

where  $k_{fY}$  and  $k_{fL}$ , the coefficients of additional mass for linear motion along the Y- and Z-axes, are presented in figure 7 as a function of the fireness ratio in the pion view and the maximum depth-to-width ratio of the equivalent ellipsoid.

The additional moment of inertia of a fuselage rotating about one of its principal axes is determined from the equations of rotational momentum for the equivalent colline id as presented in references 3 and 5. For rotations about the Y and Z axes, respectively, the additional momenta of inertia may be expressed in terms of the average width and depth of the fuselage as

$$I_{Y_{\Omega}} = \frac{\rho}{5} k^{\dagger}_{fY} L_{f} wd \left( \frac{L_{e}^{2}}{\mu} + \frac{3d^{2}}{2\pi} \right)$$
 (8)

and

$$I_{Z_{a}} = \frac{o}{5} k^{\dagger}_{fZ} L_{f} wd \left( \frac{L_{f}^{2}}{4} + \frac{3w^{2}}{2\pi} \right)$$
 (9)

where the coefficients of additional moments of inertia  $k^{\dagger}_{fY}$  and  $k^{\dagger}_{fZ}$  are presented in figure 8 as a function of the fineness ratio in the plan view and the maximum depth-to-width ratio of the equivalent ellipsoid.

For rotations about reference swinging exes displaced from but parallel, respectively, to the Y and Z principal exes of the fuselage the values of the additional moments of inertia may be obtained from equations (6) to (9) and the moment-of-inertia transference equation

$$I_a = I_a + m_a t^2$$

as

$$I_{Y_{a}}^{\dagger} = \frac{\rho}{5} k_{fY}^{\dagger} L_{fW}^{\dagger} d \left( \frac{L_{f}^{2}}{4} + \frac{3d^{2}}{2\pi} \right) + \rho \left( k_{fZ}^{\dagger} L_{fW}^{\dagger} d l_{fY}^{2} \right)$$
 (10)

and

$$I_{Z_a}^{\dagger} = \frac{\rho}{5} k^{\dagger}_{fZ} L_{f}^{Wd} \left( \frac{L_f^2}{\mu} + \frac{3w^2}{2\pi} \right) + \rho \left( k_{fY} L_f^{Wd} l_{fZ}^2 \right)$$
 (11)

For motion along and rotation about the X principal axis of the airplane, calculations and theory show that the values of the additional mass and additional moments of inertia are relatively small; accordingly, they have not been considered in the following equations for the estimation of the total additional moments of inertia of the airplane about its swinging exes.

#### Complete Airplane

The evaluation of the additional moments of inertia for a complete airplane entails the applications of equations (1) to (11) to the various portions of the airplane, which can be considered either as flat plates or ellipsoidal bodies. For each part of the airplane the value of the additional moment of inertia is computed about an axis through the centroid of area of the surface and parallel to the corresponding reference swinging axis by applying

the appropriate forms of equations (2), (3), (4), (8), and (9). The total values of the additional mements of inertia about the reference swinging axes are then obtained by using the appropriate forms of the mement-of-inertia transference formulas (equations (5), (10), and (11)).

It has been found that, inasmuch as some terms of the resulting general equation for the complete airplane are small, approximate equations for additional moments of inertia about each of the three swinging axes may be written without sacrificing accuracy. These approximate equations are:

$$I_{X_a}^{\dagger} = \frac{\pi \rho}{48} \left( k^* D_{\lambda} D_{\Gamma} S^2 b \right)_W + \rho \left( k_{f \Upsilon}^{\dagger} L_{f}^{\dagger} w d l_{f \Upsilon}^{2} \right)_{\text{fus}}$$
 (12)

$$I_{Y_a} = \left[ \frac{\rho}{5} k_{fY}^{I} L_{fWd} \left( \frac{L_f^2}{4} + \frac{3d^2}{2\pi} \right) \right]_{fus}$$

+ 
$$\rho \left(k_{fZ}L_{f}wdl_{f_{Y}}^{2}\right)_{fus}$$
 +  $\frac{\pi\rho}{4}\left(k_{g}^{2} + \frac{\kappa^{2}}{b} l_{t_{Y}}^{2}\right)_{ht}$  (13)

$$I_{Z_{\mathbf{a}}}' = \left[\frac{\rho}{5} \, \mathbf{k'}_{fZ} L_{f} \text{wd} \left(\frac{L_{f}^{2}}{4} + \frac{3w^{2}}{2\pi}\right)\right]_{fus}$$

+ 
$$\rho \left(k_{fY}L_{f}^{vd}l_{fZ}^{2}\right)_{fus}$$
 +  $\frac{\pi\rho}{4}\left(k\frac{s^{2}}{b}l_{tz}^{2}\right)_{vt}$  (14)

where, for convenience in making the calculations, S/b is substituted for  $\bar{c}$ .

An illustration of the application of the procedure to determine the additional moments of inertia is given in the appendix for a typical fighter airplane. The method of determining the various dimensions and areas is indicated on figures 1 and 2. The values of the terms left out of the approximate equations are also given and it may be seen that these terms are negligible.

Comparison of Experimental and Calculated Results

An estimate of the accuracy of the equations (12) to (14) used in calculating the additional moments of inertia may be obtained by inspection of figures 9 to 11 in which the calculated values are plotted against the experimental values which were determined by swinging the models in a vacuum chamber as described in reference 5. The agreement is good since the experimental errors may have been as high as 50 percent of the true values, which were only a small percentage of the total values of the moments of inertia measured with respect to the swinging axes.

A comparison of the experimental values of the tetal moments of inertia about the body axes with the additional moments of inertia about the axes of rotation is presented in table I. This table indicates that with a swinging gear similar to the arrangement shown in figure 1, the additional moments of inertia probably do not exceed, in the majority of cases, 25 percent of the true moments of inertia about the body axes.

#### SUMMARY OF RESULTS

Formulas have been developed for the estimation of the additional moments of inertia of airplanes or airplane models. A correlation of the experimental data on 40 models indicates that the satisfactory estimations of the values of the moments of inertia, due to the ambient-air effect, may be determined by means of these formulas.

Langley Memorial Aeronautical Laboratory
National Advisory Committee for Aeronautics
Langley Field, Va., October 15, 1946

### APPENDIX

## EXAMPLE OF METHOD FOR CALCULATING THE ADDITIONAL MOMENTS OF INERTIA ABOUT THE REFERENCE SWINGING

#### AXES FOR A TYPICAL FIGHTER AIRPLANE

Pertinent data. In order to illustrate the use of the formulas presented in the text, calculations are presented for a typical fighter airplane. Figure 1 is a sketch of the airplane in position for swinging.

Fertinent coefficients and dimensional data for the airplane are as follows:

Wing:	
Aroa, S, sq ft	250
Spen, b, ft	36
Aspect ratio, A	5.9
Mean chord, c, ft (S/b)	6.1
Component of distance in the X-Y plane from centreid of area	
of wing to Y-axis of rotation, $l_{W_Y}$ , ft	
Taper ratio, $\lambda$	
Dihedral, $\Gamma$ , dog	-
Additional mass coefficient, k (from fig. 3	-
for $A = 5.9$ )	0.95
Additurnal meant-of-inertia coefficient, kt, for the	••//
X gay, ging (from fig. 4 for $A = 9.9$ )	0.88
For the Y said sug, the reciprocal of the espect	0.0
/ 1. 1	
ratio ( ) to used to give an elaitional moment-of-	
ironola oce fice out k! from fixore 4	0.10
Toper ratio our school, $\mathcal{D}_{\lambda}$ (from Fig. 6)	0.12
Dibodial correction, $D_{\Gamma}$ (from fig. 5)	0.96
Fuselcge:	
Length, Le, ft	23.5
Average width, w, ft	3.40
Average depth, d., ft	3.8
Distance from centroid of side area of fusclage to	.,, ,
X-exis of rotation, $l_{ extbf{f}_{ extbf{X}}}$ , ft $\dots$	7.68
Component of distance in X-Y plane from controld of top	
area of fuselage to Y-axis of rotation, 1fx, ft	4.50

Fuselage - Continued:
Distance from centroid of side area of fuselage to
Z-axis of rotation, $l_{f_Z}$ , ft
Additional moment-of-inertia coefficients
obtained from the width-depth ratios and
fineness ratio of the fuselage (fig. 7)
key
k <sub>fY</sub>
•
Additional moment-of-inertia coefficients obtained
from the width-depth ratios and fineness ratio of
the fuselage (fig. 8)
k <sup>†</sup> fY · · · · · · · · · · · · · · · · · · ·
k <sup>†</sup> fY
Horizontal tail:
Area, S, sq ft
Span, b, ft
Aspect ratio, A
Meen chord, c, ft (S/b)
Component of distance in the X-Y plane from centraid of area of horizontal tail to Y-axis of
rotation, l <sub>ty</sub>
Additional mass coefficient, k, (from fig. 3)0.90
Additional moment-of-inertia coefficient, k, for
the X swinging (from fig. 4 for
$A = 3.70 \dots \dots$
For the Y awinging the reciprocal of the aspect
ratio $\left(\frac{1}{3.7}\right)$ is used to give an additional moment- of-inertia coefficient ki from figure $h$
of-inertia coefficient k! from figure h
Taper ratio, $\lambda$
Taper ratio, $\lambda$
Vertical tail:
Area, S, sq ft
Span, b, ft 4.6
Aspect ratio, A
Mean chord, c, ft (S/b)
Distance from centroid of area to X-axis of
rotation, $l_{t_X}$ , ft 4.3
Distance from centroid of area to Z-axis of
rotation, $l_{t_7}$ , ft
Additional mass coefficient, k (from fig. 3) 0.59
Additional moment-of-inertia coefficient, k, for the
X swinging (from fig. 4 for A 1.04) 0.41

Vertical tail - Continued:

For the Z swinging the reciprocal of the aspect

X-axis. The additional moment of inertia about the X swinging axis (shown in fig. 1) is obtained from the approximate equation (12) as follows:

$$I_{X_{\mathbf{a}}} = \frac{\rho \pi}{48} \left( k^* D_{\lambda} D_{\Gamma} s^2 b \right)_{W} + \rho \left( k_{fY} L_{f} w d l_{fX} \right)_{fus}$$
(A1)

Substituting the proper tabulated values in equation (Al) gives

$$I_{X_{8}}^{t} = \frac{3.14 (0.002378)}{48} [(0.88)(0.86)(0.88)(220)^{2}(36)]_{W}$$

$$+ [(0.002378)(1.04)(23.5)(3.40)(3.85)(7.68)^{2}]_{fus}$$

$$= (185.66)_{W} + (45)_{fus}$$

$$= 231 \text{ slug-ft}^{2}$$
(A2)

The value of  $I_{X_A}^{-1}$  determined experimentally was 269 slug-feet<sup>2</sup>.

The exact equation for the additional moments of inertia about the X swinging axis contains the following terms, a numerical evaluation of which gives quantities that are indicative of the relatively small magnitudes of these terms:

$$\frac{\rho\pi}{48} \left( k^{2}D_{\lambda}D_{\Gamma}s^{2}b \right)_{ht}$$

$$= \left[ \frac{3.14(0.002378)}{48} (0.78)(0.85)(1)(34.7)^{2}(11.5) \right]_{ht}$$

$$= 1.47 \text{ slug-ft}^{2}$$
(A3)

$$\frac{\rho\pi}{48} \left( k^{\bullet} D_{\lambda} S^{2} b \right)_{vt}$$

$$= \left[ \frac{3.14(0.002378)}{48} (0.41)(0.72)(20.2)^{2}(4.6) \right]_{Vt}$$

$$= 0.089 \text{ slug-ft}^{2} (Ah)$$

$$\frac{\rho\pi}{4} \left( k \frac{S^2}{b} l_{tX}^2 \right)_{vt}$$

$$= \left[ \frac{3.14(0.002378)}{4} (0.59) \frac{(20.2)^2}{4.6} (4.3)^2 \right]_{vt}$$

$$= 1.81 \text{ slug-ft}^2$$
(A=)

Y-axis. The additional moment of inertia about the Y swinging axis (shown in fig. 1) is obtained from the approximate equation (13) as follows:

$$I_{Y_{a}}^{\dagger} = \left[\frac{\rho}{5} k^{\dagger}_{fY} L_{f}^{\dagger} wd \left(\frac{L_{f}^{2}}{4} + \frac{3d^{2}}{2\pi}\right)\right]_{fus} + \left[\rho \left(k_{fZ}^{\dagger} L_{f}^{\dagger} wd l_{fY}^{2}\right)\right]_{fus}$$

$$+ \left[\frac{\rho\pi}{4} \left(k \frac{S^{2}}{b} l_{tY}^{2}\right)\right]_{ht} \qquad (A6)$$

Substituting the proper tabulated values in equation (A6) gives

$$I_{Y_{a}}^{\dagger} = \left\{ \frac{0.002378}{5} (0.89)(23.5)(3.40)(3.85) \left[ \frac{(23.5)^{2}}{4} + \frac{3}{6.28}(3.85)^{2} \right] \right\}_{fus}$$

$$+ \left[ (0.002378)(0.86)(23.5)(3.40)(3.85)(4.50)^{2} \right]_{fus}$$

$$+ \left[ \frac{3.14(0.002378)[0.90(3^{1}4.7)^{2}}{11.5} (15.8)^{2} \right]_{ht}^{2}$$

= 
$$(18.90)_{\text{fus}} + (19.74)_{\text{fus}} + (43.99)_{\text{ht}}$$
  
=  $76 \text{ slug-ft}^2$  (A7)

The value of  $I_{Y_{\alpha}}^{-1}$  determined experimentally was 143 slug-feet<sup>2</sup>.

The exact equation for the additional moments of inertia about the Y swinging axis contains the following terms, a numerical evaluation of which gives quantities that are indicative of the relatively small magnitudes of these terms:

$$\frac{\rho\pi}{48} \left( k^i D_{\chi} D_{\Gamma} S^2 \right)_{W}$$

$$= \left[\frac{3.14(0.002378)}{48}(0.12)(0.86)(0.96)(220)^{2}(6.1)\right]$$

$$= 4.56 \text{ slug-ft}^{2} \tag{A8}$$

$$\frac{\rho\pi}{4} \left( k \frac{8^2}{b} v_{\text{WY}}^2 \right)_{\text{W}} = \left\{ \frac{3.14(0.002378)}{4} \left[ \frac{0.95(220)^2}{35} (1.6)^2 \right] \right\}_{\text{W}}$$

$$= 6.11 \text{ slng-ft}^2 \tag{A9}$$

$$\frac{\rho\pi}{48} \left( k^3 D_{\lambda} D_{\underline{1}} s^2 \overline{e} \right)_{ht}$$

$$= \left[\frac{3.14(0.002378)}{48}(0.18)(0.85)(1)(34.7)^{2}(3.1)\right]_{ht}$$

$$= 0.089 \text{ slug-ft}^{2} \tag{A1c}$$

Z-axis. - The additional moment of inertia about the Z awinging axis (shown in fig. 1) is obtained from the approximate equation (14) as follows:

(A12)

$$I_{Z_{2}}' = \left[\frac{\rho}{5} k^{\dagger}_{fZ} L_{f}^{vd} \left(\frac{L_{f}^{2}}{4} + \frac{3v^{2}}{2\pi}\right)\right]_{fus} + \left[\rho \left(k_{fY} L_{f}^{vd} l_{fZ}^{2}\right)\right]_{fus} + \left[\frac{\rho\pi}{4} \left(k \frac{g^{2}}{b} l_{tZ}^{2}\right)\right]_{vt}$$
(All)

Substituting the proper tabulated values in equation (All) gives

$$I_{Z_{B}}' = \left\{ \frac{(0.002378)}{5} (0.94)(23.5)(3.4)(3.85) \left[ \frac{(23.5)^{2}}{4} + \frac{3}{6.28}(3.4)^{2} \right] \right\}_{\text{fus}}$$

$$+ \left[ (0.002378)(1.04)(23.5)(3.4)(3.85)(19.1)^{2} \right]_{\text{fus}}$$

$$+ \left\{ \frac{3.14(0.002378)}{4} \left[ \frac{0.59(20.2)^{2}}{4.6} (30.6)^{2} \right] \right\}_{\text{vt}}$$

$$= (19.70)_{\text{fus}} + (277.65)_{\text{fus}} + (91.63)_{\text{vt}}$$

$$= 389 \text{ slug-ft}^{2}$$
(A12)

The value of  $I_{Z_a}$  determined experimentally was 345 slug-feet<sup>2</sup>.

The exact equation for the additional moments of inertia about the Z swinging axis contains the following term, a numerical evaluation of which gives a quantity that is indicative of the relatively small magnitude of this term:

$$\frac{\rho\pi}{48} \left( k^4 D_{\lambda} \Omega^{2} \delta \right)_{V_{\tau}^{\pm}}$$

$$= \left[ \frac{3.15(0.402378)}{48} (0.41)(0.72)(20.2)^{2}(4.4) \right]_{V_{\tau}^{\pm}}$$

$$= 0.083 \text{ slug-St}^{2}$$
(A13)

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TABLE I. - COMPARISON OF THE EXPERIMENTAL TOTAL VALUES OF THE MOMENTS OF INERTIA ABOUT THE BODY AXES AND THE ADDITIONAL MOMENTS OF INERTIA ABOUT THE AXIS OF ROTATION OF THE COMPOUND PENDULUM OF 40 FREE-SPINNING AIRPLANE MODELS

[Values are in gr-in.2]

<del>                                     </del>		·					
Airplane model	Model scale	IX	IXa	I <sub>Y</sub>	Iya;	$\mathbf{I}_{\mathbf{Z}}$	Iza
шопол	Board		<u>~</u>				
XSB2U-1	1/16	6,520	1160	14,082	398	19,217	2088
BT-9	1/16	5,280	717	7,452	344	11,260	1270
P-35	1/18	3,195	300	4,681	160	7,095	385
XOS2U-1	1/16	4,014	877	10,190	859	13,234	1849
P-36A	1/20	1,327	232	2,937	79	3,962	482
XS02U-1	1/16	5,188	890	16,369	540 630	20,361	2337 1894
YFM-1	1/25	8,514	1128	4,902	639	12,556	
XSO3C-1	1/14	11,566	1385	34,169	1291	41,896 6,115	3309 445
XSB2A-1	1/25	1,908	246	4,417	1.50 653		1601
B-26 XP-46	1/26	11,266	523 166	12,354	651 116	22,899	395
XTBU-1	1/20 1/24	2,158	350	3,640 6,328	119	5,617 9,216	579
XP -47B	1/20	3,311 9,111	320 387	8,572	244	16,978	656
XP -67	1/27	6,156	183	3,752	63	9,327	415
XSB3C 1	1/20.83	8,694	831	11,232	729	19,008	826
XP -69	1/20	17,375	1181	32,307	735	48,451	2153
XP -62	1/22	5,402	432	9,198	233	13,755	843
XF14C-1	1/20	7,695	695	9,686	368	15,990	1694
XF2A-1	1/16	4,578	436	7,452	503	11,250	875
Λ-1.7	1/15	17,190	2100	21,265	730	34,464	2233
XF4F-2	1/12	12,730	2540	31,149	1725	41,006	2740
NF - 1	1/12	16,289	1700	32,882	1308	48,098	5630
XF4U~L	1/20	4,862	383	5,303	268	9,474	779
XF5F-1	1./22	4,407	299	2,927	102	7,044	313
XBT-12	1/14	9,744	1320	16,305	601	24,606	1603
XBT-11	1/16	5,416	876	8,746	335	11,835	1202
XP -50	1/25	2,965	98	1,630	99	4,560	122
XTBF-1	1/22	4,793	857	8,632	386	12,722	958
P-39D	1/20	3,417	137	3,992	90	7,032	432
P-30	1/16	7,944	738	14,511	1.80	20,301	1550
P-40F	1/20	3,304	274	5,190	88	7,980	1246
SNC -1 XSB 2D-1	1/14 1/20	4,856	659 566	11,194 16,775	49 <u>1</u> 463	15,394 24,856	1601
1		9,155		1	11.56	30,136	2195
XA-39 XP-75	1/20	13,915 20,309	1529 1745	17,542 33,701	1089	51,829	3814
XP -71	1/28	18,930	653	12,860	370	31,790	3241
DC -3	1/23.75	18,554	2764	25,517	<b>870</b>	41,856	2461
XCG-16	1/32	2.821	707	2,043	150	4,688	752
XC-82	1/60	770	1.67	373	144	1,119	45
SBN-1	1/18	3,587	428	6,601	239	9,741	867
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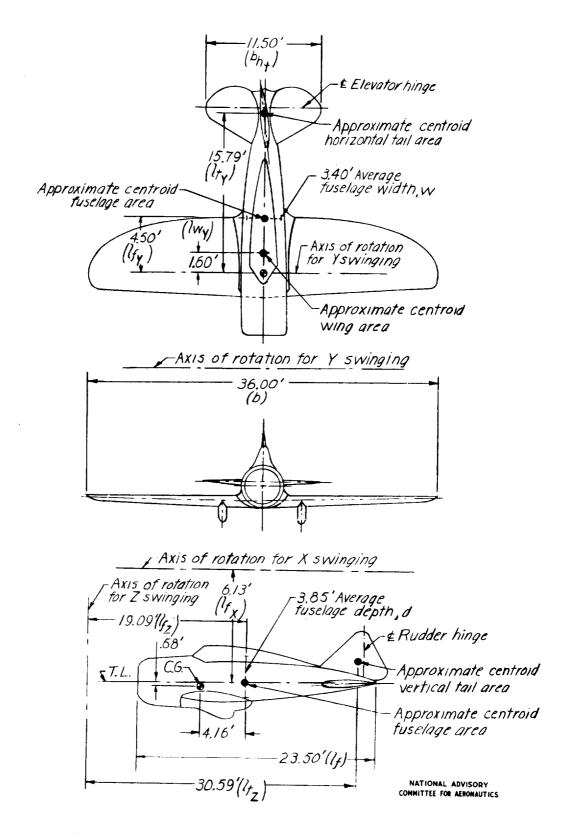
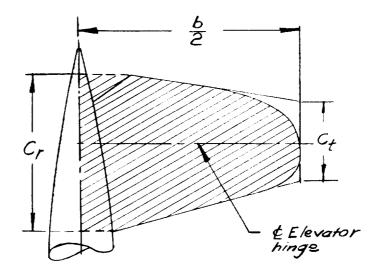


Figure 1.- Sketch of a typical model showing dimensions necessary for the calculation of the additional moments of inertia. All dimensions are full-scale values.



(a) Horizontal tail.

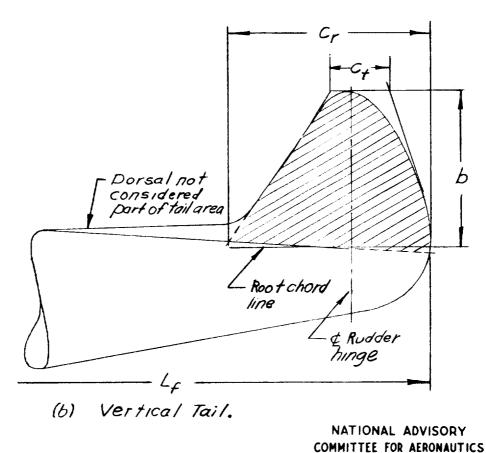


Figure 2.- Sketch illustrating method of measuring the chords, span, and area of the horizontal and vertical tail surfaces;

 $c_t = (\frac{2S}{b}) - c_r$  where S is the area of the tail surfaces.

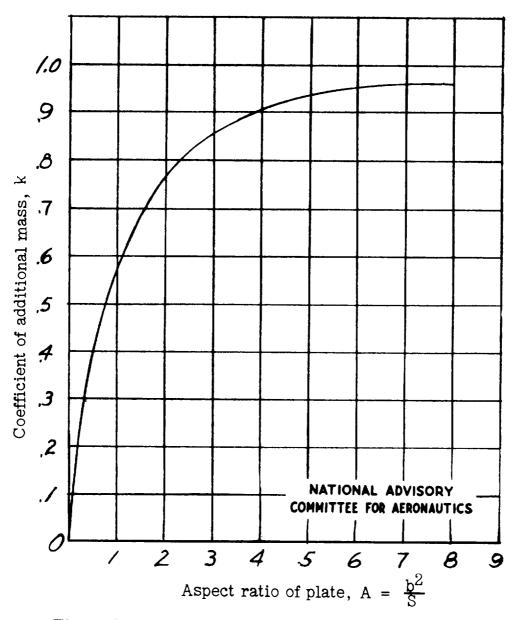
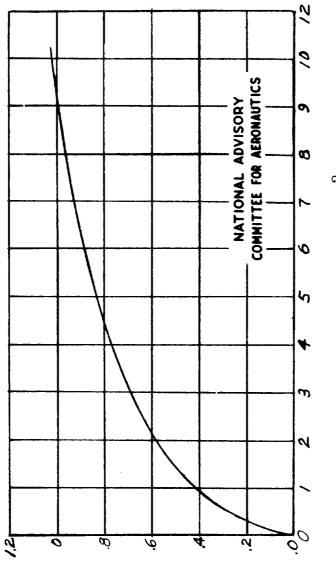


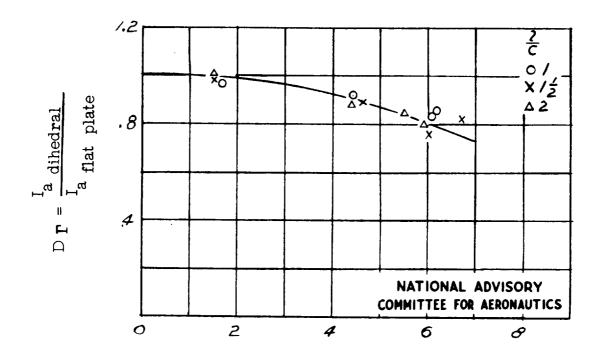
Figure 3.- Coefficients of additional mass for rectangular plates. Curve is average of NACA 1933 and 1940 tests. (Reference 5.)



Coefficient of additional moment of inertia, k'

Curve is average of NACA 1933 and 1940 tests. (Reference 5.) Coefficients of additional moment of inertia for rectangular Figure 4.plates.

Aspect ratio of plate, A



Dihedral angle, deg

Figure 5.- Variation of the additional moments of inertia of a single plate with dihedral angle; A = 4 (reference 5).

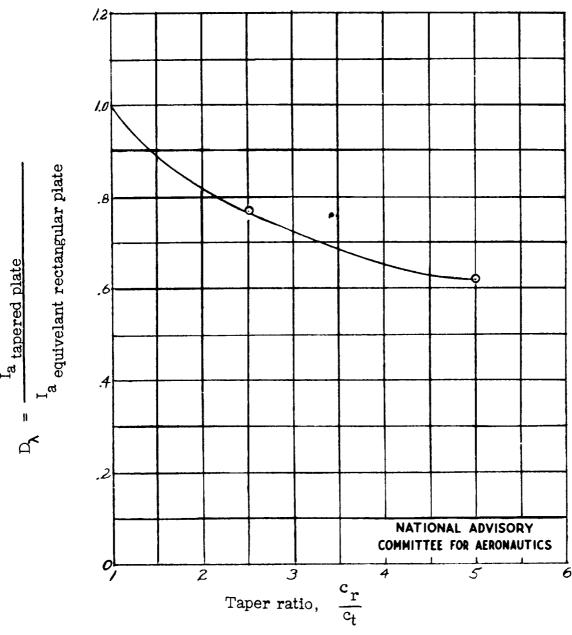


Figure 6.- Dependence of the additional moment of inertia on taper ratio (reference 5).

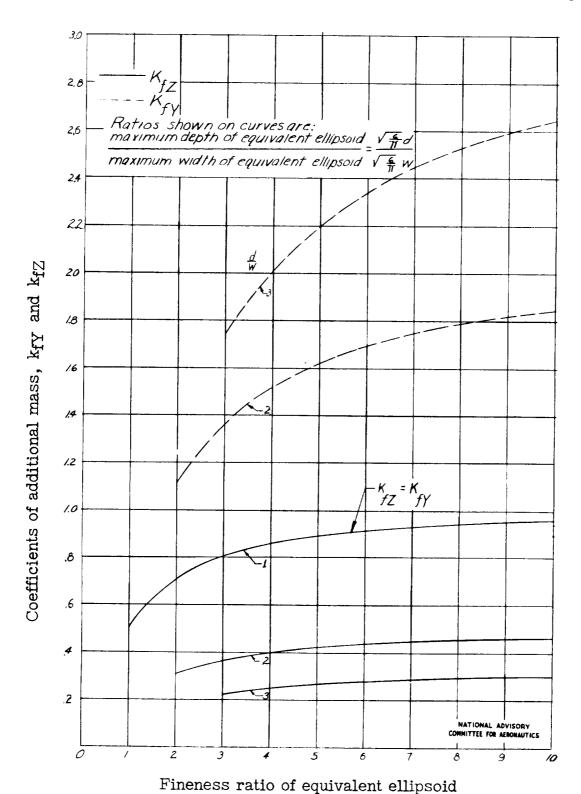


Figure 7.- The variation of the coefficients of additional mass  ${\bf k}_{fY}$  and  ${\bf k}_{fZ}$  with the fineness ratio of an equivalent ellipsoid

$$\frac{\text{Actual length}}{\text{Maximum width}} = \frac{L_{f}}{w\sqrt{\frac{6}{\pi}}}.$$

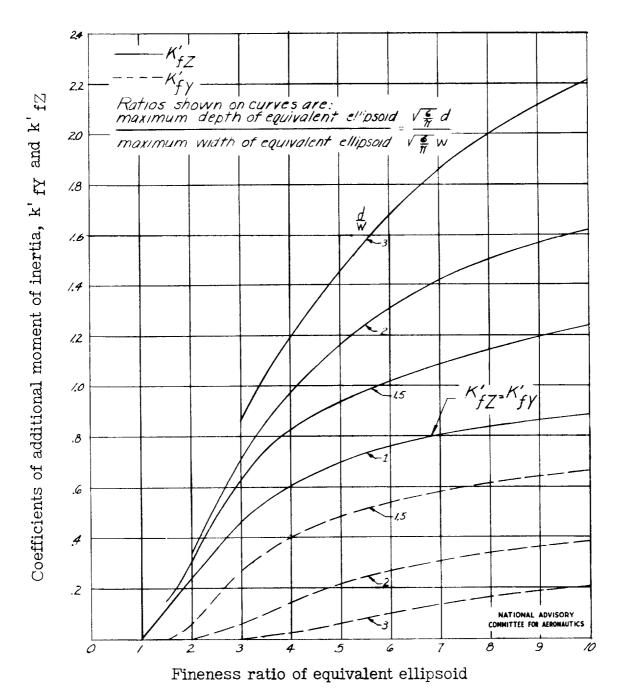


Figure 8.- The variation of the coefficients of additional moment of inertia k'  $_{\rm fZ}$  and k'  $_{\rm fY}$  with fineness ratio of an equivalent

ellipsoid 
$$\frac{\text{Actual length}}{\text{Maximum width}} = \frac{L_f}{w\sqrt{\frac{6}{\pi}}}$$

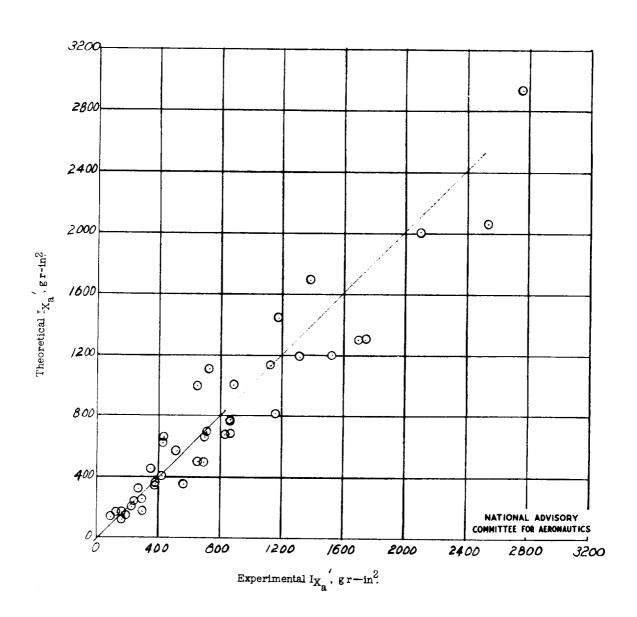


Figure 9.- Comparison of the experimental and calculated model values of the additional moments of inertia about the X swinging axis for 40 free-spinning models.

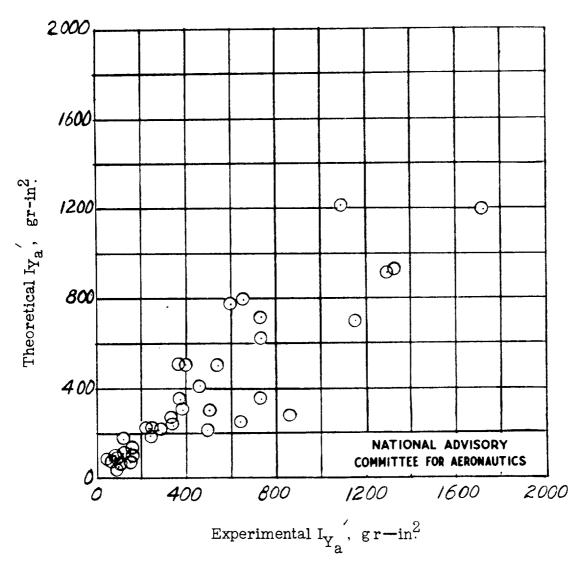


Figure 10.- Comparison of the experimental and calculated model values of the additional moments of inertia about the Y-swinging axis for 40 free-spinning models.

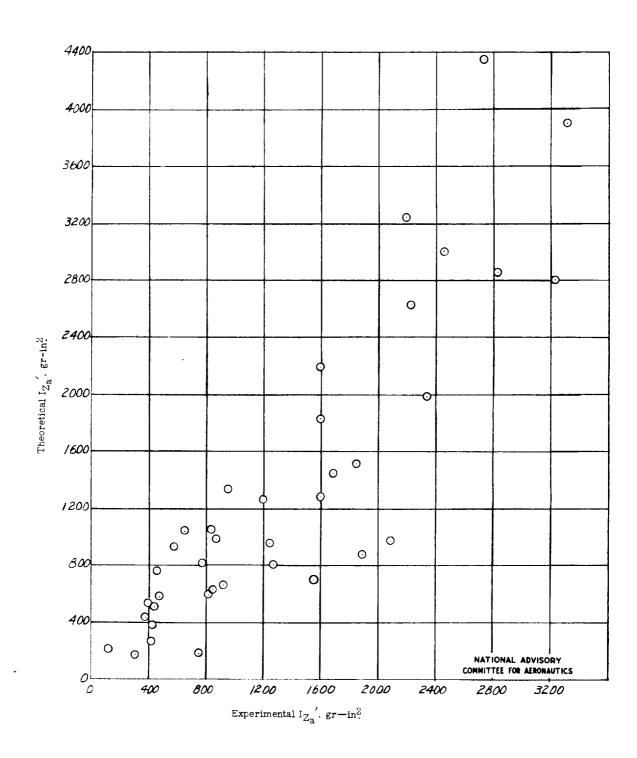


Figure 11.- Comparison of the experimental and calculated model values of the additional moments of inertia about the Z swinging axis for 40 free-spinning models.